

Potential theory of Dirichlet forms with jump kernels blowing up at the boundary

Panki Kim

Seoul National University

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References

This talk is based on a joint project with Renming Song (University of Illinois, USA) and Zoran Vondraček (University of Zagreb, Croatia).

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[KSV] K, Song & Vondraček, Potential theory of Dirichlet forms with jump kernels blowing up at the boundary, arXiv:2208.09192[math.PR].

- 1 **Motivation by Two examples**
 - Example 1: Trace process
 - Example 2: Non-local Neumann problems
- 2 **Estimates of resurrection kernel in general case**
- 3 **Dirichlet forms with jump kernels blowing up at the boundary**
 - Setup
 - Decay rate of harmonic function
- 4 **Main results**

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Trace process

Let $d \geq 1$, $\mathbb{R}_+^d := \{x = (\tilde{x}, x_d) : x_d > 0\}$ and $\mathbb{R}_-^d := \{x = (\tilde{x}, x_d) : x_d < 0\}$.

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$$\Delta^{\frac{\alpha}{2}} f(x) = \text{p.v.} \int_{\mathbb{R}^d} (f(y) - f(x)) C_{d,\alpha} |x - y|^{-d-\alpha} dy,$$

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Define $A_t := \int_0^t \mathbf{1}_{\mathbb{R}_+^d}(X_s) ds$ and let $\tau_t := \inf\{s > 0 : A_s > t\}$ be its right-continuous inverse.

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The process $Y = (Y_t)_{t \geq 0}$, defined by $Y_t = X_{\tau_t}$ called the **trace** process of X on \mathbb{R}_+^d .

Trace process through resurrection

The part of the trace process Y until its first hitting time of the boundary $\partial\mathbb{R}_+^d = \{(\tilde{x}, 0) : \tilde{x} \in \mathbb{R}^{d-1}\}$ can be described in the following way:

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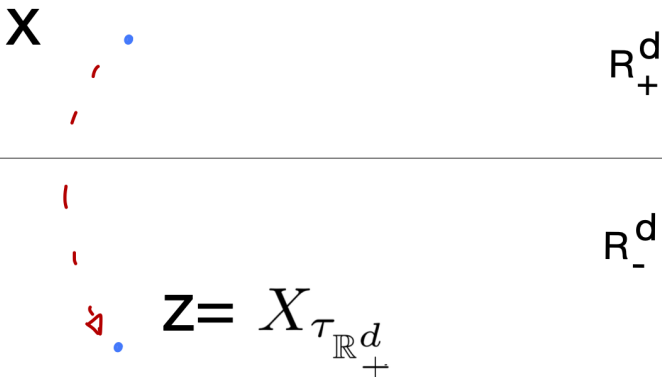
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$$\mathbf{X} = X_{\tau_{\mathbb{R}_+^d}-}$$

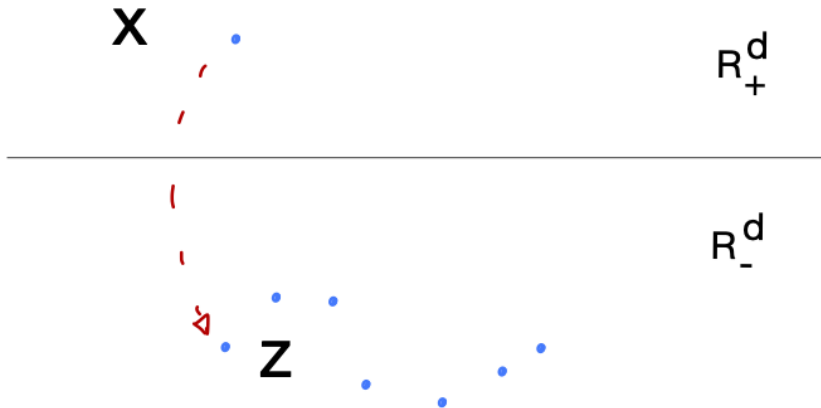

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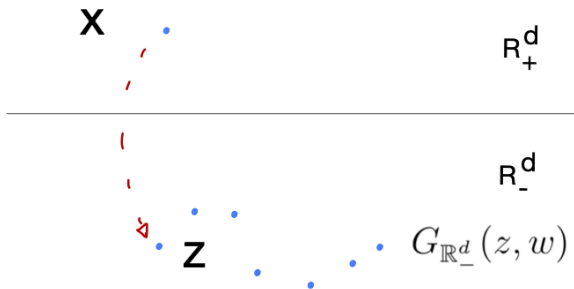
$G_{\mathbb{R}_-^d}^X(z, w)$, the **Green function** of the process X killed upon exiting \mathbb{R}_-^d . That is,

$$\mathbb{E}_z \int_0^{\tau_{\mathbb{R}_-^d}} \mathbf{1}_A(X_t) dt = \int_A G_{\mathbb{R}_-^d}^X(z, w) dw.$$

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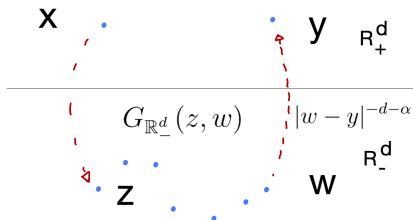
For $z \in \mathbb{R}_-^d$, $\mathbb{P}_z(X_{\tau_{\mathbb{R}_-^d}} \in dy) = P_{\mathbb{R}_-^d}(z, y) dy = \mathbb{E}_z \int_0^\tau |X_s - y|^{-d-\alpha} ds dy$ on \mathbb{R}_+^d .

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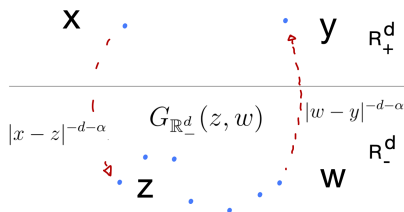
When X jumps out of \mathbb{R}_+^d from the point x , we continue the process by resurrecting it at $y \in \mathbb{R}_+^d$ according to the kernel

$$q(x, y) := \int_{\mathbb{R}_-^d} |x - z|^{-d-\alpha} P_{\mathbb{R}_-^d}(z, y) dz, \quad x, y \in \mathbb{R}_+^d. \quad (1.1)$$

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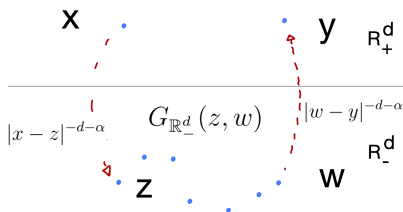


We will call $q(x, y)$ a **resurrection kernel**.

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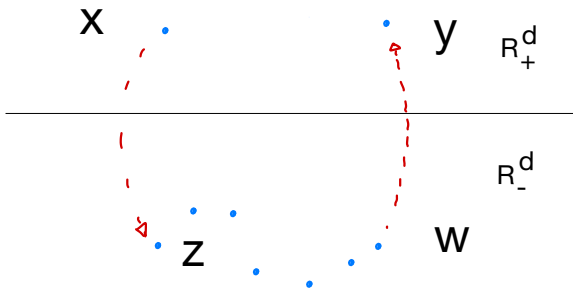
Since $G_{\mathbb{R}_-^d}^X(\cdot, \cdot)$ is symmetric, it follows that $q(x, y) = q(y, x)$ for all $x, y \in \mathbb{R}_+^d$.

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It follows from [Bogdan, Grzywny, Pietruska-Paluba & Rutkowski, 20] that (in case $d \geq 3$),

$$J(x, y) \asymp q(x, y) \asymp |x - y|^{-d-\alpha} \left(\frac{|x - y|^2}{x_d y_d} \right)^{\alpha/2}, \quad x_d \wedge y_d \leq |x - y|.$$

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This asymptotic relation shows that the jump kernel $J(x, y)$ **blows up** with rate $x_d^{-\alpha/2}$ when x approaches the boundary $\partial\mathbb{R}_+^d$.

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Non-local Neumann problems with corresponding resurrection kernel

Another motivation is the process introduced in

- Dipierro, Ros-Oton & Valdinoci. Nonlocal problems with Neumann boundary conditions. *Rev. Mat. Iberoam.* 33 (2017), 377-416.
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For that process, the resurrection kernel $q(x, y)$ is given by (1.1) with the Poisson kernel $P_{\mathbb{R}^d_-}(z, y)$ replaced by $|z - y|^{-d-\alpha} / \int_{\mathbb{R}^d_+} |z - w|^{-d-\alpha} dw$.

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That is,

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The jump kernel of this process **blows up** with rate $\log |x_d|$ when x approaches the boundary $\partial\mathbb{R}_+^d$.

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Resurrection kernel in general case

We generalize these two examples by replacing the Poisson kernel $P_{\mathbb{R}_+^d}(z, y)$ and the kernel $|z - y|^{-d-\alpha} / \int_{\mathbb{R}_+^d} |z - w|^{-d-\alpha} dw$ by a **very general return kernel** $p(z, y)$.

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Let the kernel $y \mapsto p(z, y)$, $z \in \mathbb{R}_-^d$, $y \in \mathbb{R}_+^d$, be a **probability density** such that the corresponding resurrection kernel

$$q(x, y) = \int_{\mathbb{R}_-^d} |x - z|^{-d-\alpha} p(z, y) dy, \quad x, y \in \mathbb{R}_+^d,$$

is **symmetric**.

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This flexibility in choosing the return kernel allows us to obtain resurrection kernels with **various blow-up rates** at the boundary.

A general upper bound

$$q(x, y) \leq c(x_d \wedge y_d)^{-d-\alpha}$$

An example of general forms

We consider the following general resurrection kernel:

$$q(x, y) := c \int_{\mathbb{R}_-^d} \Psi \left(\frac{|y - z|^2}{y_d |z_d|} \right) \frac{|z_d|^\alpha}{|y - z|^{d+\alpha}} \frac{dz}{|x - z|^{d+\alpha}}, \quad x, y \in \mathbb{R}_+^d,$$

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Then

$$q(\lambda x, \lambda y) = \lambda^{-d-\alpha} q(x, y)$$

and

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Moreover,

$$q(x, y) = q(y, x).$$

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(a) For the trace process of an isotopic α -stable process on \mathbb{R}_+^d ,

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is the Poisson kernel for \mathbb{R}_-^d .

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is the Poisson kernel for \mathbb{R}_-^d . Thus, here $\Psi(r) = cr^{\alpha/2}$.

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$$q(x, y) := c \int_{\mathbb{R}_+^d} \Psi \left(\frac{|y - z|^2}{y_d |z_d|} \right) \frac{|z_d|^\alpha}{|y - z|^{d+\alpha}} \frac{dz}{|x - z|^{d+\alpha}}, \quad x, y \in \mathbb{R}_+^d,$$

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$p(z, y) = |z - y|^{-d-\alpha} / \mu(z)$ where $\mu(z) = \int_{\mathbb{R}_+^d} |z - y|^{-d-\alpha} dy$. Since

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As a corollary of this general result, we have the following.

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Corollary: Let $\gamma \in (-\infty, 1 \wedge \alpha)$ and $\delta \in \mathbb{R}$. Suppose $\Psi(t) = t^\gamma \log^\delta t$, $t \geq 2$,

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(2) For $x, y \in \mathbb{R}_+^d$ with $x_d \wedge y_d \leq |x - y|$, it holds that

$$q(x, y) \asymp J(x, y) \asymp |x - y|^{-d-\alpha} \begin{cases} \left(\frac{|x-y|^2}{x_d y_d} \right)^\gamma \log^\delta \left(\frac{|x-y|^2}{x_d y_d} \right) & \text{when } \gamma > 0; \\ \log^{\delta+1} \left(\frac{|x-y|^2}{x_d y_d} \right) & \text{when } \delta > -1, \gamma = 0; \\ \log \left(e + \log \left(\frac{|x-y|^2}{x_d y_d} \right) \right) & \text{when } \delta = -1, \gamma = 0; \\ 1 & \text{when } \delta < -1, \gamma = 0 \text{ or } \gamma < 0. \end{cases}$$

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Let $d \geq 1$, $\alpha \in (0, 2)$ and assume that $0 \leq \beta_1 \leq \beta_2 < 1 \wedge \alpha$.

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$$C_1(R/r)^{\beta_1} \leq \frac{\Phi(R)}{\Phi(r)} \leq C_2(R/r)^{\beta_2}, \quad 2 \leq r < R < \infty.$$

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(A2) If $\alpha \geq 1$, there exists $\theta > \alpha - 1$ such that for every $a > 0$ there exists $C = C(a) > 0$ such that

$$|\mathcal{B}(x, y) - \mathcal{B}(x, x)| \leq C \left(\frac{|x - y|}{x_d \wedge y_d} \right)^\theta \quad \text{for all } x, y \in \mathbb{R}_+^d \text{ with } x_d \wedge y_d \geq a|x - y|.$$

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Remark: (1) **(A3)** $\Rightarrow \mathcal{B}(x, y) \geq c_1 > 0$. (2) **(A4)** $\Rightarrow \mathcal{B}(x, x) \equiv c_2 > 0$.

Dirichlet form with critical killing (Potential) term

For $\kappa \in [0, \infty)$, set

$$\mathcal{E}^\kappa(u, v) := \frac{1}{2} \int_{\mathbb{R}_+^d} \int_{\mathbb{R}_+^d} (u(x) - u(y))(v(x) - v(y)) J(x, y) dy dx + \int_{\mathbb{R}_+^d} u(x)v(x) \kappa x_d^{-\alpha} dx.$$

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Under assumptions **(A1)**-**(A4)**, for all $\kappa \in [0, \infty)$, there exists a symmetric, scale invariant and horizontally translation invariant Hunt process

$Y^\kappa = ((Y_t^\kappa)_{t \geq 0}, (\mathbb{P}_x)_{x \in \mathbb{R}_+^d})$ associated with $(\mathcal{E}^\kappa, \mathcal{F}^\kappa)$.

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Decay rate and the constant on the killing function

We now associate the constant κ from the killing function $x_d^{-\alpha}$ with a positive parameter $p = p_\kappa > 0$ which will be the **decay rate of harmonic functions**.

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$$C(\alpha, q, \mathcal{B}) = \begin{cases} \int_{\mathbb{R}^{d-1}} \int_0^1 \frac{(s^q - 1)(1 - s^{\alpha-q-1})}{(1-s)^{1+\alpha}} \frac{\mathcal{B}((1-s)\tilde{u}, 1), s\mathbf{e}_d}{(|\tilde{u}|^2 + 1)^{(d+\alpha)/2}} ds d\tilde{u} & \text{if } d \geq 2 \\ \int_0^1 \frac{(s^q - 1)(1 - s^{\alpha-q-1})}{(1-s)^{1+\alpha}} \mathcal{B}(1, s) ds & \text{if } d = 1. \end{cases}$$

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Then $C(\alpha, 0, \mathcal{B}) = C(\alpha, \alpha - 1, \mathcal{B}) = 0$ and the function $q \mapsto C(\alpha, q, \mathcal{B})$ is strictly increasing and continuous on $[(\alpha - 1)_+, \alpha - \beta_2)$.

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When $\Phi(r) = r^\beta$ with $\beta \in (0, 1 \wedge \alpha)$, it holds that $\lim_{q \uparrow \alpha - \beta} C(\alpha, q, \mathcal{B}) = \infty$, so $\kappa \mapsto p_\kappa$ is an increasing bijection from $[0, \infty)$ onto $[(\alpha - 1)_+, \alpha - \beta)$.

In the remainder of this talk we will fix $\kappa \in [0, \lim_{q \uparrow \alpha - \beta_2} C(\alpha, q, \mathcal{B})]$, and assume $\alpha > 1$ if $\kappa = 0$ so that $p_0 = \alpha - 1 > 0$.

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$$L^{\mathcal{B}} f(x) = L^{\mathcal{B}, \kappa} f(x) = \text{p.v.} \int_{\mathbb{R}_+^d} (f(y) - f(x)) J(x, y) dy - C(\alpha, p, \mathcal{B}) x_d^{-\alpha} f(x), \quad x \in \mathbb{R}_+^d,$$

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Hence the operator $L^{\mathcal{B}}$ annihilates the p -th power of the distance to the boundary.

Outline

- 1 **Motivation by Two examples**
 - Example 1: Trace process
 - Example 2: Non-local Neumann problems
- 2 **Estimates of resurrection kernel in general case**
- 3 **Dirichlet forms with jump kernels blowing up at the boundary**
 - Setup
 - Decay rate of harmonic function
- 4 **Main results**

Scale invariant boundary Harnack principle with exact decay rate

Theorem 2

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Assume that \mathcal{B} satisfies **(A1)**-**(A4)** and $\kappa = C(\alpha, p, \mathcal{B})$. Then there exists $C \geq 1$ such that for all $r > 0$, $\tilde{w} \in \mathbb{R}^{d-1}$, and any non-negative function f in \mathbb{R}_+^d which is harmonic in $B((\tilde{w}, 0), 2r) \cap \mathbb{R}_+^d$ with respect to Y^κ and vanishes continuously on $B((\tilde{w}, 0), 2r) \cap \partial\mathbb{R}_+^d$, we have

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$$\frac{f(x)}{x_d^p} \leq C \frac{f(y)}{y_d^p}, \quad x, y \in B((\tilde{w}, 0), r/2) \cap \mathbb{R}_+^d.$$

Sharp two-sided estimates for the Green function

Theorem 3

Suppose that $p \in (0, \alpha - \beta_2) \cap [(\alpha - 1)_+, \alpha - \beta_2)$ and that \mathcal{B} satisfies **(A1)**-**(A4)** and $\kappa = C(\alpha, p, \mathcal{B})$.

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$G^\kappa : \mathbb{R}_+^d \times \mathbb{R}_+^d \rightarrow [0, \infty]$ such that $G^\kappa(x, \cdot)$ is continuous in $\mathbb{R}_+^d \setminus \{x\}$. Moreover, $G^\kappa(x, y)$ has the following estimates: for all $x, y \in \mathbb{R}_+^d$,

$$G^\kappa(x, y) \asymp \begin{cases} \left(\frac{x_d}{|x-y|} \wedge 1 \right)^p \left(\frac{y_d}{|x-y|} \wedge 1 \right)^p \frac{1}{|x-y|^{d-\alpha}}, & \alpha < d; \\ \left(\frac{x \wedge y}{|x-y|} \wedge 1 \right)^p \log \left(e + \frac{x \vee y}{|x-y|} \right), & \alpha = 1 = d; \\ \left(\frac{x \wedge y}{|x-y|} \wedge 1 \right)^p (x \vee y \vee |x-y|)^{\alpha-1}, & \alpha > 1 = d. \end{cases}$$

The results in this talk can be considered as a counterpart of our previous work on jump kernels vanishing at the boundary:

- [KSV1] K, Song & Vondraček: On potential theory of Markov processes with jump kernels decaying at the boundary. *Potential Analysis* (2023).
- [KSV2] K, Song & Vondraček, Sharp two-sided Green function estimates for Dirichlet forms degenerate at the boundary. To appear in *Journal of the European Mathematical Society (JEMS)*, 2023.
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The research in [KSV1]-[KSV3] was motivated by the fact that subordinate killed Lévy processes have jump kernels vanishing at the boundary.

Combined with jump kernel vanishing at the boundary

As a particular case,

$$\mathcal{B}(x, y) \asymp \left(\frac{x_d}{|x - y|} \wedge 1 \right)^\beta \left(\frac{y_d}{|x - y|} \wedge 1 \right)^\beta \asymp \left(\frac{x_d y_d}{|x - y|^2} \wedge 1 \right)^\beta = \tilde{B}(x, y).$$

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In the decay case with \mathcal{B} as above, we have already proved the BHP and the Green function estimate with the decay rate equal to the p -th power of the distance to the boundary. Thus, we now have these results for all $-(1 \wedge \alpha) < \beta < \infty$.

No killing case

Suppose $\kappa = 0$, $\alpha > 1$, **(A2)** and **(A4)** hold and

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The corresponding non-local operator is

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This covers the case of the censored α -stable process ($\mathcal{B}(x, y) \equiv 1$), $\alpha \in (1, 2)$, studied in Bogdan, Burdzy & Chen, 03, in the half-space case.

Thank you.
감사합니다.